The effect of Mn on mineral stability in metapelites revisited:

new \(a-x\) relations for manganese-bearing minerals

R. W. White\(^1\)*, R. Powell\(^2\) & T. E. Johnson\(^1\)#

\(^1\)Institute of Geoscience, University of Mainz, D-55099 Mainz, Germany
\(^2\)School of Earth Sciences, University of Melbourne, Vic 3010, Australia
\(^3\)Department of Earth Sciences, University of Cambridge, Cambridge, CB2 3EQ, UK

# Present address: Department of Applied Geology, The Institute for Geoscience Research (TIGeR), Curtin University, GPO Box U1987, Perth, WA 6845, Australia

*Corresponding author: rwhite@uni-mainz.de

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ABSTRACT

The $a$-$x$ relations recently presented in White et al. 2014, *Journal of Metamorphic Geology*, 32, 261–286 are extended to include MnO. This provides a set of internally consistent $a$-$x$ relations for metapelitic rocks in the MnO–Na$_2$O–CaO–K$_2$O–FeO–MgO–Al$_2$O$_3$–SiO$_2$–H$_2$O–TiO$_2$–O$_2$ (MnNCKFMASHTO) system. The mixing parameters for the Mn-bearing minerals were estimated using the new micro-$\phi$ approach of Powell et al. 2014, *Journal of Metamorphic Geology*, 32, 245–260. Then the Mn-end-member thermodynamic properties were calibrated using a database of co-existing minerals involving literature data from rocks and from experiments on natural materials. Mn-end-members were calibrated for orthopyroxene, cordierite, staurolite, chloritoid, chlorite, biotite, ilmenite and hematite, assuming known properties for the garnet end-member spessartine. The addition of MnO to phase diagram calculations results in a marked expansion of the stability of garnet-bearing assemblages. At greenschist facies conditions garnet stability is extended down temperature. At amphibolite facies conditions the garnet-in boundary shifts to lower pressure. While the addition of MnO greatly influences the stability of garnet, it has relatively little effect on the stability of other common metapelitic minerals, with the resultant diagrams being topologically very similar to those calculated without MnO. Furthermore, the addition of MnO in the amounts measured in most metapelites has only a small effect on the mode of garnet, with calculated garnet modes remaining smaller than 1% in the $P$–$T$ range outside its predicted Mn-free $P$–$T$ range.
INTRODUCTION

While the key equilibria that control the metamorphic mineral assemblages in metapelites can be shown graphically in systems as simple as K$_2$O–FeO–MgO–Al$_2$O$_3$–SiO$_2$–H$_2$O (KFMASH), the quantitative analysis of such rocks using phase equilibria requires larger systems (e.g. White et al., 2000, 2007, 2014; Diener et al., 2007; Smye et al., 2010). These larger systems are necessary to accommodate relatively minor components such as TiO$_2$, Na$_2$O, CaO Fe$_2$O$_3$ and MnO that nevertheless may exhibit a strong influence over the stability of certain common minerals. For example, TiO$_2$ is known to strongly affect the stability of biotite at higher temperature.

Manganese is well known to affect mineral assemblages in metasedimentary rocks via its strong stabilising effect on garnet (e.g. Atherton, 1964; Osberg, 1971; Symmes & Ferry, 1992; Droop & Harte, 1995; Mahar et al., 1997; Tinkham et al., 2001). Thus, despite its typically low concentration (≈ 0.1–0.3 wt %) in metasediments (e.g. Ague, 1991; Atherton & Brotherton, 1982) manganese forms a critical component in phase equilibria and the interpretation of metamorphic assemblages (e.g. Symmes & Ferry, 1992; Droop & Harte, 1995; Mahar et al., 1997). Calculations in manganese-free systems ranging from KFMASH to NCKFMASHTO (Na$_2$O–CaO–K$_2$O–FeO–MgO–Al$_2$O$_3$–SiO$_2$–H$_2$O–TiO$_2$–O$_2$) show a relatively restricted stability of garnet, even in relatively Fe-rich bulk compositions (e.g. Powell et al., 1998; White et al., 2000, 2014). For example, the NCKFMASHTO calculations presented in White et al. (2014) for the amphibolite facies average metapelite composition from Ague (1991) have garnet stability restricted to pressures above ≈7.5 kbar and temperatures above ≈570 °C. These calculations contrast markedly with the common occurrence of manganese-bearing garnet in many greenschist and amphibolite facies rocks, including those from contact aureoles and low-P Buchan type terrains (e.g. Harte & Hudson, 1979; Hudson, 1985; Pattison & Tinkham, 2009)

Some of the first studies to undertake quantitative phase diagram calculations beyond the KFMASH system (e.g. Spear & Cheney, 1989; Symmes & Ferry, 1992; Mahar et al., 1997) incorporated Mn-end-members in the calculations in order to better predict the stability of garnet. While a large number of studies since have considered MnO in phase diagram calculations in various systems (e.g. Tinkham & Ghent, 2005; White et al., 2005; Wei et al., 2007; Pattison & Tinkham, 2009), these studies generally did not involve the
development of new thermodynamic models or recalibration of the Mn end-member data. For example the $a$–$x$ relations presented in Tinkham et al. (2001) and Tinkham & Ghent (2005), are simply the combination of the Mahar et al. (1997) MnO relationships to the extant Mn-free $a$–$x$ relations of the time (e.g. Holland et al., 1998; Holland & Powell, 1998; Powell & Holland 1999; White et al., 2000). Obviously this is not a defensible approach given that the underlying formulation of the Mahar et al. (1997) $a$–$x$ relationships was ideal mixing, while the later models with which they were combined involved the symmetric formalism. Thus, there were large inherent inconsistencies in the $a$–$x$ relations used in these previous studies.

In this paper we develop new end-member properties for the MnO-bearing end-members following a similar approach to Mahar et al. (1997). This builds on and extends the thermodynamic descriptions of White et al. (2014) and they are compatible with the most recent Holland & Powell (2011) dataset. Thus, for the first time since 1997, the thermodynamic properties of the manganese end-members for phases are consistent with the other end-members in each phase.

THE THERMODYNAMICS OF MN-BEARING MINERALS

The following phases, with their phase abbreviations, are considered to be Mn-bearing: garnet (g), orthopyroxene (opx), cordierite (cd), staurolite (st), chlorite (chl), chloritoid (ctd), biotite (bi), and ilmenite (ilm). The manganese end-members of these minerals are the same as those whose properties appear in ds6 (Holland & Powell, 2011), except for the manganese end-member of orthopyroxene which is made from the dataset end-member, pyroxamangite (MnSiO$_3$): garnet (spessartine, spss, Mn$_3$Al$_2$Si$_3$O$_{12}$), biotite (mnbi, K Mn$_3$Al$_2$Si$_3$O$_{10}$(OH)$_2$, orthopyroxene (mnopx, Mn$_2$Si$_2$O$_6$), cordierite (mncrd, Mn$_2$Al$_4$Si$_5$O$_{18}$), staurolite (mnst, Mn$_4$Al$_{18}$Si$_7$O$_{44}$(OH)$_4$), chloritoid (mnctd, MnAl$_2$SiO$_5$(OH)$_2$), chlorite (mnchl, Mn$_5$Al$_2$Si$_3$O$_{10}$(OH)$_8$) and ilmenite (pyrophanite, pnt, MnTiO$_3$).

In biotite, orthopyroxene and chlorite, manganese is distributed across non-equivalent octahedral sites such that the site fraction of Mn is the same on the different sites (coined equidistribution). Note that this is not equipartition, which in this case might mean making Mn/Fe$^{2+}$ the same across the sites. Unlike equipartition, equidistribution is not
thermodynamically inconsistent (Holland & Powell, 2007). While it is likely that Mn is partitioned between the non-equivalent sites as Mg and Fe$^{2+}$ are, rather than equally distributed, such an approach would require additional ordered manganese end-members, and the enthalpy of ordering would need to be estimated. Although approaches to implement this are available, e.g. Powell et al. (2014), given the lack of data on Mn ordering and the fact that Mn is nearly always a minor constituent of these phases, the additional complexity of such an approach is not considered to be warranted.

The formulation of the $a-x$ relations of the minerals involves ideal-mixing-on-sites and the symmetric formalism (Powell & Holland, 1993; Powell et al., 2014). The parameterisation follows White et al. (2014) closely, with the addition of the non-ideality involving Mn substitution in the minerals remaining to be considered. Except for garnet and ilmenite, there are no experimental data that would allow the pairwise macroscopic interaction energies ($W$) between the manganese and the other end-members in the minerals to be estimated. As argued in Powell et al. (2014), the simplification of assuming that all the interaction energies, $W$, involving the manganese and other end-members are zero is unlikely to be correct, and is likely to bias the results of calculations. Given that even Mg–Fe$^{2+}$ interactions have significant non-zero $W$, it is much better to give the $a-x$ relations involving Mn a good “shape”, with non-zero $W$. A way to do this is presented in Powell et al. (2014), referred to as micro-$\phi$ (see below), and this approach has already been used to generate most of the $W$ involving ferrous end-members in the minerals in White et al. (2014). Once the shape is established, what remains to be determined is the enthalpy adjustment ($\Delta H$, or DQF in THERMOCALC parlance) of the manganese end-members already in the dataset (Holland & Powell, 2011), which involves the “historical” values relating to the ideal-mixing-on-sites approach of Mahar et al. (1997).

In micro-$\phi$, the idea is to implement an expectation that the $a-x$ relations for the $X-Y$ subsystem of a mineral be related to the those of the Mg–Al subsystem, commonly the subsystem that is best known experimentally. In this $X$ might be, for example, Fe$^{2+}$ or Mn, and $Y$, Fe$^{3+}$ or Cr. Here, the microscopic $w$ of the $X-Y$ subsystem is made proportional to the Mg–Al subsystem with the proportionality represented by $\phi$. This allows the $X-Y$ subsystem to be easily made less ($\phi < 1$) or more ($\phi > 1$) non-ideal than the Mg–Al subsystem. The implementation of micro-$\phi$ starts with disassembling the macroscopic $W$ into their constituent microscopic same-site and cross-site $w$ that arise
from pairwise interactions between the cations on the sites in the mineral (Powell & Holland, 1993). Then simplifications and approximations are made to these $w$ using heuristics\(^1\), before reassembling the macroscopic $W$. In Powell et al. (2014) and White et al. (2014), with $X = \text{Fe}^{2+}$ and $Y = \text{Fe}^{3+}$, the main heuristic values allowing parameterisation of the $a$-$x$ relations are:

$$
\begin{align*}
\phi &= 0.7 \\
\phi_3 &= 0.8 \\
w_{\text{FeMg,oct}} &= 4 \text{ kJ} \\
w_{\text{MgAl,oct}} &= 10 \text{ kJ},
\end{align*}
$$

(1)

in which $\phi$ has the effect, $w_{\text{FeX,oct}} = \phi w_{\text{MgX,oct}}$, where Mg, Fe\(^{2+}\), and $X$ are mixing on the octahedral site, oct. Thus, the Fe subsystem of a mineral can be made to have non-ideality that is proportional to that of the Mg subsystem with the proportionality arising from $\phi$. Similarly, $\phi_3$ has the effect, $w_{\text{YFe3,oct}} = \phi_3 w_{\text{YAl,oct}}$, the Fe\(^{3+}\) subsystem being made to have non-ideality that is proportional to that of the Al subsystem. In this way, for example,

$$w_{\text{FeFe3,oct}} = \phi \phi_3 w_{\text{MgAl,oct}}.$$  

Making $w_{\text{MgMn,oct}} \approx 2.9 \text{ kJ}$, or rounded to 3 kJ for phases where little is known about the overall $a$-$x$ relations of the mineral, keeping the $W$ as whole numbers. Here, as opposed to the above, the value for $w_{\text{FeMn,oct}}$ is taken as known and the micro-$\phi$ approach is used to derive the value for $w_{\text{MgMn,oct}}$. Regarding the behaviour of Mn in mixing with Fe\(^{2+}\) and Mg, the data is very limited and no data exists for directly constraining Mn mixing properties involving octahedral sites in silicates. However, there is some information regarding garnet and ilmenite, through experiments on Fe–Mn exchange experiments (Pownceby et al., 1987; Feenstra & Peters, 1996). As recognised by Pownceby et al. (1987), the sundry correlations

\(^1\)Heuristic is used here in the sense of an experience-based rule of thumb or educated guess, commonly in the context of providing a numerical value for a parameter or a relationship between parameters
between controlling thermodynamic parameters means that separating information about
garnet from that of ilmenite via the experimental data is difficult. However, Pownceby et
al. (1987) do argue that Fe–Mn interactions have to be slightly stronger in ilmenite than in
garnet: they suggest on a single site basis $W_{\rm FeMn}^{\text{ilm}} - \frac{1}{3} W_{\text{FeMn}}^{\text{g}} \approx 1.2 \text{ kJ}$. Separately, O’Neill
(1998) estimated from Fe–Mn exchange between ilmenite and olivine that
$W_{\text{FeMn}}^{\text{ilm}} = 1.8 \pm 0.1 \text{ kJ}$, with provisos about the thermodynamics of olivine. This value is
also consistent with O’Neill et al. (1989), who derived $W_{\text{FeMn}}^{\text{ilm}} = 2.2 \pm 0.3 \text{ kJ}$. This leads to
our adoption of $W_{\text{FeMn}}^{\text{ilm}} = 2 \text{ kJ}$. This then leads to $W_{\text{alm,spss}} = 3(W_{\text{FeMn}}^{\text{ilm}} - 1.2) \approx 2 \text{ kJ}.$

Recently, Dachs et al. (2014) undertook a detailed analysis of Fe–Mn mixing in garnet
and suggested almandine-spessartine mixing is asymmetric and a little more non-ideal than
that adopted here, for more almandine-rich garnet. In the light of the discussion of the
various experimental data by O’Neill (1998), adoption of the simpler, symmetric, model is
defensible. Making the unavoidable step of carrying this value across to octahedral sites in
silicates, leads to $w_{MgMn,\text{oct}} = 3 \text{ kJ}$, once $\phi_{\text{Mn}}$ is chosen to be 0.7, given that this $w$ is likely
to be of the same order as $w_{MgFe,\text{oct}}$. The interaction energies that result from the adoption
of these heuristics are given in Appendix 1. They should be seen in the context of the
Appendix in White et al. (2014) for the way in which the approach of Powell & Holland
(1993) and Holland & Powell (1996ab) for writing $a$-$x$ relations is implemented.

The dataset of Holland & Powell (2011) contains data for all the manganese
end-members above but their enthalpies are based on the original ideal-mixing calculations
of Mahar et al. (1997). As these are superceded by the non-ideal mixing calculations
presented here, the enthalpies of these end-members need to be modified. A reference point
is provided by the properties of the pyrophanite and spessartine dataset end-members
(Holland & Powell, 2011), as they are based on experimental data rather than being from
Mahar et al. (1997). They provide an anchor for establishing manganese end-member
properties for the other minerals. As in Mahar et al. (1997), the approach taken is to use
Mg–Mn exchange reactions between the minerals using a natural assemblage database
constructed for the purpose (Appendix 2). The database used for calibration is presented
as Supplementary Information. Exchange reactions are good for calibration purposes as the
resulting thermodynamic properties are not sensitive to the chosen $P$–$T$ of the natural
assemblages. This is the converse of saying that exchange reactions make bad thermometers
because they are so sensitive to the thermodynamics. Such thermometers can easily lead to
strong bias in thermometric results (e.g. through poor thermodynamic formulation or
when ferric is ignored), e.g. Powell & Holland (2008). The exchange reactions used are the
Mg–Mn ones so as to minimise the consequences of difficulties with respect to ferric iron.
The derivation of the enthalpy modifications of the dataset via the Mn–Mg exchange
reactions is given in Appendix 3, completing the thermodynamic descriptions of the phases.

**PHASE DIAGRAM CALCULATIONS**

The phase diagrams presented here using the new $a-x$ relations were calculated using
THERMOCALC version 3.37 and the internally-consistent end-member dataset of Holland &
Powell (2011), ds62 (created 6th February 2011). Calculations were undertaken in a range
of chemical systems ranging from MnO–K$_2$O–FeO–MgO–Al$_2$O$_3$–SiO$_2$–H$_2$O (MnKFMASH)
to MnO–Na$_2$O–CaO–K$_2$O–FeO–MgO–Al$_2$O$_3$–SiO$_2$–H$_2$O–TiO$_2$–O$_2$ (MnNCKFMASHTO).
The version of the Holland & Powell dataset used here is slightly newer than that used in
White et al. (2014) for modelling in NCKFMASHTO (ds61), though the differences in the
resulting diagrams from this are very minor.

In addition to the Mn-bearing minerals, garnet, orthopyroxene, cordierite, staurolite,
chlorite, chloritoid, biotite and ilmenite, the following minerals were considered in the
phase diagram calculations: sillimanite/kyanite/andalusite (sill/ky/and), spinel (sp),
muscovite (mu), paragonite (pa), margarite (ma), K-feldspar (ksp), plagioclase (pl), albite
(ab), epidote (ep), sphene (sph), magnetite (mt), rutile (ru), quartz (q) and melt (liq). For
the minerals not considered to be Mn-bearing, the $a-x$ relationships of White et al. (2014)
were used for the white micas and silicate melt; Holland & Powell (2003) for feldspar,
White et al. (2000) for magnetite at subsolidus conditions, White et al. (2002) for
magnetite and spinel at suprasolidus conditions, and Holland & Powell (2011) for epidote.

**MnKFMASH**

**P–T Projections and compatibility diagrams**

In adding MnO to the KFMASH system, each of the invariant points in KFMASH becomes
a univariant line extending away from the KFMASH invariant, typically to lower $P-T$
conditions. The Mn content of the phases increase away from the invariant point. Where
the KFMASH invariant point is garnet bearing, the univariant that emanates from it may extend a considerable distance in \( P-T \). By contrast, MnKFMASH univariant reactions that lack garnet typically only extend a short distance from their KFMASH origins before garnet appears at a new MnKFMASH invariant. A MnKFMASH grid is shown in Fig. 1 for subsolidus (450 °C) to suprasolidus (950 °C) conditions. The Mn content of garnet \((m(g))\) along four of the univariant reactions is shown via the horizontal ticks along each reaction. The grid is relatively simple with only two invariant points. The suprasolidus univariant reactions extend only a short distance from their KFMASH origins before terminating in another subsystem. The first invariant point \([\text{opx cd ksp liq}]\)—using the \([\text{absent phase}]\) notation—occurs at \(\approx 0.8 \text{kbar and } \approx 480 \text{ °C}\) and links the main univariant reactions seen in the subsolidus part of the grid. Details of the reactions that emanate from this point are shown in the oval shaped inset in Fig. 1. The second invariant point \([\text{opx st chl ctd}]\) occurs where the muscovite breakdown reaction intersects the wet solidus (\(\approx 5.8 \text{kbar, } \approx 725 \text{ °C}\)). For \(P-T\) conditions below the solidus and the muscovite breakdown reaction there are four univariant reactions that could be seen under typical crustal metamorphic conditions. These reactions, in order of increasing \(T\) are:

\[
\begin{align*}
\text{als + ctd} &= \text{g + st + chl}, \\
\text{chl + ctd} &= \text{g + st + bi}, \\
\text{st + chl} &= \text{g + als + bi} \\
\text{g + als + chl} &= \text{cd + bi},
\end{align*}
\]

where als represents the stable aluminosilicate of and, sill and ky. A singularity occurs along reaction 6 involving als swapping sides of the reaction at lower pressure. For reactions 3–5, which all emanate from KFMASH invariant points, each represents a garnet-bearing equivalent to the garnet-absent reaction from each KFMASH invariant (White et al., 2014), with the KFMASH and MNKFMASH reactions typically occuring only a few degrees apart for most of their length because the incorporation of Mn is minor.
in all of the phases apart from garnet. As a consequence, the garnet-absent assemblage possible for each MnKFMASH reaction (e.g. ky–st–ctd–chl for below reaction 3) typically has a very limited $P$–$T$ stability range between the KFMASH and MNKFMASH reactions.

The phase relationships relating to the low variance MnKFMASH equilibria (with $v = 1$–3) can also be shown in compatibility diagrams, that additionally show the composition space that these equilibria occupy. A series of AFM and MnFM compatibility triangles are shown in Fig. 2 and Fig. 3 respectively for the $P$–$T$ conditions shown on Fig. 1 as small triangles. The AFM triangles (Fig. 2) have garnet, quartz, muscovite and H$_2$O in excess and the MnFM diagrams (Fig. 3) have aluminosilicate, quartz, muscovite and H$_2$O in excess.

The AFM compatibility diagrams are shown for an isobaric transect at 4.5 kbar (Fig. 2a–f) at the $P$–$T$ conditions shown as open triangles in Fig. 1. As garnet is taken as in-excess, these diagrams show the changing stable divariant to quadrivariant equilibria amongst the remaining ferromagnesian phases and aluminosilicate. The sequence of AFM diagrams shows the changes in stable tie-triangles on crossing each of the univariant reactions 3–5. For example Fig. 2a and Fig. 2b illustrate the breakdown of the ky–ctd tie-line and the formation of the st–chl tie line on crossing reaction 3, with the phases involved having very similar compositions to those expected in the MnO-free system.

The mineral composition relationships either side of reaction 3 are shown in a sequence of MnFM compatibility triangles (Fig. 3) calculated for pressures of 12 kbar and 6 kbar (filled triangles on Fig 1). These diagrams show the varying size of the garnet one-phase field as a function of pressure along the reaction. As these diagrams are calculated with kyanite in excess, they are only appropriate for rather aluminous compositions. For the diagrams at 12 kbar (Fig. 3a, b), the garnet one-phase field is relatively large and extends a considerable distance from the MnO apex towards the FeO apex. Chlorite, staurolite and, where present, chloritoid all plot at low MnO contents. Thus, rocks with only small amounts of MnO may contain garnet at these conditions. At $T$ below the reaction, two stable divariant assemblages occur (g–st–ctd and g–chl–ctd). As discussed above, the garnet-absent tie triangle (st–ctd–chl) is not stable at these conditions, but appears from the base of the compatibility triangle less than 0.1 °C below the MnKFMASH univariant. Above the reaction only one stable tie triangle exists, and, as kyanite is in excess, this reaction is a terminal chloritoid reaction.
The diagrams at 6 kbar (Fig. 3c, d) show substantially reduced garnet one-phase fields in comparison with those at 12 kbar and the other phases (especially staurolite and chloritoid) may incorporate substantially more MnO before garnet must become stable. Thus, garnet is limited to somewhat higher bulk rock MnO contents and is considerably more MnO rich than at higher \( P \).

**Pseudosections**

The MnKFMASH system is useful for demonstrating the effect of adding MnO on the stability of key assemblages. Figure 4 shows the predicted assemblages for three different bulk compositions with different \( x_{Fe} = \frac{FeO}{FeO + MgO} \) and \( x_{Al} = \frac{Al_2O_3}{Al_2O_3 + FeO + MgO} \) proportions. Figure 4a is a \( P-T \) pseudosection from 450 °C to 950 °C and for the same composition as used for fig. 4 in White et al. (2014) but with 0.1 mol. % MnO. Pseudosections are also presented for subsolidus conditions for a more magnesian composition (Fig. 4b) and a more magnesian and more aluminous composition (Fig. 4c), each with an MnO content of \( \approx 0.1 \) mol. % (see Table 1 for exact compositions used). The two lower \( x_{Al} \) pseudosections (Fig. 4a, b) are topologically similar, containing largely the same assemblage fields.

The pseudosections are dominated by divariant to quadrivariant fields. However, a short segment of the MnKFMASH reaction 5 is seen in each pseudosection. This reaction controls the disposition of the main divariant fields seen in each diagram. Each of these divariant fields is equivalent to a univariant equilibria stable in the KFMASH system and are little displaced in \( P-T \) from the KFMASH univariants. For example the divariant g–st–chl–bi field in each pseudosection is equivalent to the KFMASH univariant reaction:

\[
g + chl = st + bi. \tag{7}\]

In Fig. 4, the garnet-in lines for a bulk rock MnO content of 0.05 (dashed red line labelled MnO=0.05) and for the MnO-free system (red line labelled MnO=0) are also shown. These garnet-in lines illustrate the relationship between the bulk MnO content and the \( P-T \) stability of garnet, with the garnet-in lines moving down \( T \) and \( P \) relative to the MnO-absent garnet-in lines. At lower \( T \) (greenschist facies) the garnet-in boundaries are relatively steep such that the main effect of adding MnO is a shift to lower \( T \) of the garnet-in line. Under amphibolite facies conditions, the garnet-in line is somewhat flatter.
such that the main effect of adding MnO is to increase garnet stability down $P$. The
combination of these two features results in the greatest displacement of the garnet-in line
occurring at close to 500 $^\circ$C where it forms a distinct wedge that may extend to the base of
the diagram. Thus, at low $P$ the calculations predict the appearance of garnet soon
followed by its disappearance. However, in reality, such a narrow field could conceivably be
crossed with no garnet growth having occurred.

A noticeable feature of all three pseudosections is the small triangular field of
garnet-absent assemblages within the overall garnet stability field. These garnet-absent
fields each emanate from the divariant g–st–chl–bi fields in Figs 4a–c and reflect the fact
that garnet is a reactant in reaction 7. The high-$T$ boundary of this garnet absent field is
defined by the equivalent to the garnet producing KFMASH reaction,

$$st = g + als + bi.$$ (8)

Thus, along a prograde path, the calculations predict the appearance, disappearance and
reappearance of garnet.

While the stability of subsolidus garnet-bearing assemblages is strongly influenced by
MnO, the high temperature assemblages are significantly less so. At temperatures above
about 750 $^\circ$C the addition of 0.1 mol. % MnO to the bulk rock composition stabilises
garnet down pressure by less than 1 kbar (Fig. 4a). For upper amphibolite facies
conditions, manganese can have significant effect on garnet stability, particularly at 3–5
kbar where the presence of MnO may substantially extend g–sill–bi and
g–sill–cd–bi-bearing assemblages to lower $P–T$.

**MnNCKFMASHTO**

The addition of MnO to the NCKFMASHTO calculations presented in White *et al.* (2014)
results in the MnNCKFMASHTO system that closely approximates the composition space
of natural metapelites. Thus, this system is potentially useful for applied phase equilibria
modelling studies.

Figure 5 is a $P–T$ pseudosection based on the composition used for fig. 5b presented in
White *et al.* (2014), but with 0.1 mol % MnO. Garnet-in lines are additionally shown for
the MnO-free system and for MnO contents of 0.05 and 0.15. As with the MnKFMASH
calculations, the addition of MnO has a profound affect on the stability of garnet with
garnet-bearing assemblages stabilised to lower P and T. Otherwise, Fig. 5 is very similar
to the equivalent pseudosection in White et al. (2014), with the same main assemblages
present. Garnet-bearing assemblages are limited to pressures above about 7 kbar and
temperatures above about 550 °C for the MnO-free system but stabilised to pressures as
low as 2 kbar and temperatures close to 500 °C with 0.1 mol % MnO. As with the
MnKFMASH calculations, the main effect of adding MnO is to lower the temperature of
garnet stability in greenschist-facies assemblages and to lower the pressure of garnet
stability in amphibolite-facies assemblages. At pressures above 2 kbar, garnet first appears
in a series of g–chl–bi-bearing fields that may additionally involve combinations of
paragonite, epidote, plagioclase, ilmenite and magnetite. At higher temperature, the lower
pressure limit of garnet-bearing assemblages involves g–and–bi- and g–sill-bi-bearing
assemblages, in contrast to its restriction to ky-bearing assemblages in the MnO-free
system. A small, triangular garnet-absent field is also present in Fig. 5 related to the
higher variance equivalent to reaction 7.

Figure 6 shows the phase relations for a more aluminous composition, based on that of
fig. 7 in White et al. (2014). As with the pseudosection in fig. 7 of White et al. (2014),
Fig. 6 contains a series of chloritoid-bearing fields, a restricted stability range for biotite
and an enhanced stability field for the aluminosilicates compared with Fig. 5. At pressures
above about 4 kbar, the first appearance of garnet occurs in a series of biotite-absent
assemblages involving combinations of garnet, chlorite, staurolite and chloritoid. Under
greenschist-facies conditions the effect of MnO on garnet stability is less profound than in
the less aluminous composition, with the garnet-in line moving down T by about 30 °C
with the addition of 0.1 mol. % MnO. A large embayment in the garnet-in boundary exists
at close to 4 kbar, with the garnet-in line trending from about 540 °C at ≈ 3.8 kbar up to
about 640 °C at ≈ 5.8 kbar before trending back down P and T to about 510 °C and 1.6
kbar. This embayment can also be seen in the garnet-in line for the bulk MnO content of
0.05 mol. %, shown as a dashed red line. Like with the triangular garnet-absent fields
described above, this embayment is controlled by the higher variance equivalents of
reaction 7 consuming garnet and reaction 8 producing it.

Ague (1991) compiled two average metapelite compositions, one for greenschist facies
metapelites and one for amphibolite facies metapelites. Although these two compositions
are broadly similar, there are notable differences in composition between the two, including
in the bulk MnO content (Ague, 1991). In comparison with the compositions used above, the average metapelite compositions in Ague (1991) have higher CaO and Na₂O contents, but are otherwise similar to the composition used for Fig. 5, albeit with a lower $x_{\text{Fe}}$.

Figure 7 is a $P$–$T$ pseudosection for the low-grade composition from Ague (1991). A small adjustment to the CaO content was made to account for the likely presence of apatite and the $\text{Fe}_2\text{O}_3/(\text{FeO} + \text{Fe}_2\text{O}_3)$ ratio was arbitrarily set at 0.077. For Fig. 7, the MnO content (0.284 mol. %) is higher than that used in the preceding diagrams (0.1) so that most of the diagram is garnet-bearing. Garnet-absent assemblages are restricted to a small window in the bottom left of the diagram and a second at low pressure at $T > 500^\circ C$. The garnet-in lines for bulk MnO contents of 0.15 and 0.05 mol. % are additionally shown on the diagram. These two garnet-in lines are similar in shape and position to the equivalent ones in Fig. 5. Topologically, the disposition of the main AFM ferromagnesian assemblages in Fig. 7 is similar to that in Fig. 5, with the exception of the extent of garnet stability and, due to the more magnesian composition for Fig. 7, the appearance of kyanite-bearing fields in Fig. 7. As with the NCKFMASHTO calculations presented in White et al. (2014), the large number of potential phases stable in these diagrams allows for the possibility of univariant reactions to be stable. In Fig. 7 short segments of two univariant reactions are seen at conditions close to 6 kbar and 540–580 $^\circ C$ and are linked by narrow divariant fields. Each of these univariant reactions involve nine phases (excluding in-excess phases) and are dominated by large reaction coefficients for the micas and plagioclase and relatively small coefficients for the ferromagnesian phases and the oxide phases, with the exception of biotite. For example the higher temperature reaction of the two at 5.6 kbar is,

\[
44g + 1315\text{bi} + 3645\text{pa} + 1363\text{ma} = 227\text{st} + 540\text{chl} + 3617\text{pl} + 92\text{ilm} + 10\text{mt}.
\] (9)

Given the issues outlined in White et al. (2014) regarding margarite stability it is possible that the presence of these two univariants in the pseudosection is anomalous, though they may be stable in a grid.

A $P$–$T$ pseudosection for the amphibolite-facies metapelite composition from Ague (1991) is shown in Fig. 8. Here, the $\text{Fe}_2\text{O}_3/(\text{FeO} + \text{Fe}_2\text{O}_3)$ ratio was set at 0.081 but the CaO content was not reduced. The MnO content (MnO = 0.175) is lower than that for Fig. 7, resulting in a more restricted field for garnet stability. The two univariant reactions present in Fig. 7 are also seen in this bulk composition and both diagrams are topologically very similar in terms of the main assemblage fields. As with many of the other
pseudosections, a small garnet-absent field occurs within the larger area of garnet presence. Garnet-in lines are additionally shown for bulk MnO contents of 0, 0.05, 0.1 and 0.15 mol %. As with the diagrams presented earlier, the addition of MnO drives the garnet-in line down $T$ under greenschist facies conditions and down $P$ under amphibolite facies conditions, with a wedge-shaped pressure minima for garnet at approximately 2 kbar and 520 °C involving the assemblage g–and–chl–bi–pl–ilm–mt. In the MnO free system the stability of garnet-bearing assemblages is restricted to pressures above about 7.5 kbar (see White et al. (2014, fig. 11) for the corresponding NCKFMASHTO pseudosection).

The effect of MnO at higher metamorphic grades is shown on a $P$–$T$ pseudosection (Fig. 9) calculated for the amphibolite-facies metapelite composition from Ague (1991). Under amphibolite-facies conditions MnO exerts a significant effect on the stability of garnet-bearing assemblages. At temperatures below the muscovite breakdown reaction, the addition of 0.175 mol. % MnO shifts the garnet in line down approximately 3 kbar relative to the MnO-absent system. Between the muscovite breakdown reaction and the first appearance of cordierite, the garnet-in line is steep and displaced close to 80 °C down $T$ relative to the manganese-free system. Overall this greatly expands the $P$–$T$ range of g–sill–bi-bearing assemblages relative to that in the MnO-free system. At temperatures above the appearance of cordierite the effect of adding MnO is less profound, with, for example, the garnet in line now located less than 1 kbar lower than for the MnO-free system at 850 °C.

The effect of considering MnO in calculations can be further addressed via $T$–$x_{\text{MnO}}$ and $P$–$x_{\text{MnO}}$ pseudosections such as Fig. 10, based on the greenschist-facies composition from Ague (1991). Figure 10a is a $T$–$x_{\text{MnO}}$ pseudosection constructed for a pressure of 6 kbar to illustrate the down-temperature shift of the garnet-in line as a function of MnO content. The x axis varies from MnO = 0 to MnO = 0.3 mol %. At low MnO contents ($x < 0.15$) garnet is absent from the assemblages over the whole temperature range considered. The garnet-in line generally trends to lower $T$ with increasing MnO, but with several switch-backs, especially at temperatures above 590 °C, a feature that can also be seen in the garnet-in line for MnO = 0.05 in Fig. 7.

A $P$–$x_{\text{MnO}}$ pseudosection calculated for a temperature of 580 °C for the same composition range is shown in Fig. 10b and shows the down-pressure shift of the garnet-in line as a function of bulk rock MnO content. In the MnO-free system, garnet-bearing
assemblages are restricted to pressures of above 8.5 kbar for this composition but get as low as 1.5 kbar for MnO contents above 0.27 mol %. A switch-back in the garnet-in line occurs at close to 4 kbar related to the up pressure consumption of garnet through the higher-variance equivalent of reaction 8. In both the \( T-x_{\text{MnO}} \) and \( P-x_{\text{MnO}} \) pseudosections the \( P \) and \( T \) conditions for the main assemblages are little influenced by the consideration of MnO with most assemblages forming near horizontal bands across each diagram.

The pseudosections in Fig. 10 are also contoured for garnet mode. The garnet mode contours in both diagrams broadly parallel the garnet-in boundary reflecting the variable consumption or production of garnet as \( P \) or \( T \) conditions evolve. For Fig. 10a garnet is not stable for low MnO contents, such that the garnet-bearing assemblages in the diagram are a direct consequence of the addition of MnO. Despite the large increase in the overall stability of garnet-bearing assemblages, the addition of MnO results in only small proportions of garnet being stable (\(< 3\%\)) for the highest bulk MnO contents considered here. The maximum garnet contents are achieved at temperatures just below 590 °C at the low-\( T \) boundary (st-out) of the assemblage g–st–bi–pa–ma–pl–ilm–mt. For temperatures below this, relatively little garnet is produced for close to 50–100 °C above the garnet-in line. For example, at \( x = 0.5 \) garnet proportions only reach 1% approximately 100 °C above the initial appearance of garnet.

For the \( P-x_{\text{MnO}} \) pseudosection Fig. 10b the garnet mode contours are similarly sub parallel to the garnet-in line. However, unlike the \( T-x_{\text{MnO}} \) pseudosection garnet does become stable in the MnO-free system at about 8.5 kbar. As with the \( T-x_{\text{MnO}} \) diagram, the mode of garnet remains low for a considerable pressure above the garnet in line. For example at \( x = 0.7 \) the calculations predict less than 1% garnet for 3 kbar above the initial appearance of garnet. However, at higher MnO contents garnet modes above 4% are possible. In general, the mode of garnet increases with increasing pressure with the exception of the narrow field (g–sill–st–bi–pl–ilm–mt) across which garnet is consumed. A notable increase in the mode of garnet occurs across the narrow set of divariant fields and the section of univariant equilibria (at \( x < 0.75 \)) close to 5.6 kbar.
DISCUSSION AND CONCLUSIONS

The extension of the NCKFMASHTO $a-x$ relations presented in White et al. (2014) to include MnO provides a chemical system for phase diagram calculations (MnNCKFMASHTO) that closely matches that of natural metapelites and metapsammites. Importantly, the $a-x$ relations in both MnO-bearing and MnO-absent systems are thermodynamically consistent, with the Mn-end-member properties calibrated in concert with the $a-x$ relations presented in White et al. (2014). This contrasts with many previous sets of MnO-bearing $a-x$ relations in which the Mn-end-member properties (DQF adjustments) from Mahar et al. (1997) were coupled with the extant version of the Holland & Powell end-member dataset and NCKFMASHTO $a-x$ relations of the time. This resulted in inherent inconsistencies within the models, with such inconsistencies becoming more problematic as the NCKFMASHTO $a-x$ relations evolved. The thermodynamic descriptions of the Mn-bearing minerals are built on the $a-x$ development approach referred to as micro-$\phi$ of Powell et al. (2014), as detailed in the body of the paper and in Appendix 1. This approach is designed for situations where little is known about the $a-x$ relations, which is true for Mn incorporation in the minerals being considered here, apart from garnet and ilmenite. Once the $a-x$ relations were established the enthalpy modifications to the dataset properties of Holland & Powell (2011) were derived from a large database of natural coexisting minerals in metapelites (see Appendix 2–3).

The components of the MnNCKFMASHTO system commonly account for more than 98% of the mass of common metapelitic to metapsammitic rocks. Despite being a relatively minor component of metapelites, MnO exerts an important influence on the $P-T$ stability of garnet-bearing assemblages. Other than the stabilisation of garnet, the presence of MnO has little effect on the $P-T$ conditions of common assemblages. At higher grades, the effect of MnO on the $P-T$ extent of garnet-bearing assemblages is less profound, as most metapelite compositions will be garnet-bearing regardless at such conditions. Thus, for many high-$T$ studies, inclusion of MnO is likely to have little effect on the resulting estimates of $P-T$ conditions, unless the appearance or absence of garnet is of central importance. However, for modelling of transitional amphibolite to granulite facies assemblages or high-grade metamorphism at low pressures, consideration of MnO may strongly influence the interpretation of $P-T$ conditions.
Although MnO greatly expands garnet stability to lower $P-T$ conditions, the resulting mode of garnet commonly remains low until higher variance equivalents to garnet-producing KFMASH reactions are crossed. This can be seen in Fig. 10 which is contoured for garnet mode. Thus, garnet-rich rocks metamorphosed at $P-T$ conditions outside the MnO-free stability field of garnet are likely to be rather rich in MnO compared to typical metapelites.

The phase diagrams presented here reproduce the main assemblages seen in common metapelites, at least for typical MnO contents $\text{MnO} < 0.3$ wt % (e.g. Ague 1991; Atherton and Brotherton, 1982). Furthermore, using the amphibolite facies average metapelite composition from Ague (1991) as an example (Fig. 8), key features such as the garnet, staurolite, kyanite and sillimanite isograds occur in the correct order along typical metamorphic field gradients inferred for barrovian metamorphism, with the biotite isograd occurring off the diagram at $T = 400–450$ °C. For higher MnO contents, such as in Fig. 7, the garnet and biotite isograds could potentially swap positions but the other isograds would be unaffected.

In application, these models may be less reliable for bulk rock compositions much richer in MnO. Furthermore, as Mn can occur in several oxidation states in rocks, compositions rich in $\text{Mn}_2\text{O}_3$ for example are not suitable for calculations with these models. As with Fe, it is likely that all rocks contain some mixture of MnO and $\text{Mn}_2\text{O}_3$, and successful modelling of many rocks may require that small adjustments be made to the bulk rock composition. Such adjustments should be petrographically and geochemically justified, via identification of Mn in phases in which it is likely to be in the $\text{Mn}^{3+}$ state (e.g. $\text{Mn}^{3+}$ in andalusite, epidote/piedmontite) where possible.

Additionally, MnO is not considered in several key high $T$ phases (e.g. melt, sapphirine, osumilite) nor has the veracity of the Mn-end-member and mixing properties been tested at these conditions. Calculated phase equilibria at higher $T$ conditions could be erroneous, at least until Mn is incorporated in the thermodynamic descriptions of these phases.

Despite the considerable progress in the development of $a-x$ relations for complex multi-component minerals over the last ten years (e.g. White et al., 2007; Green et al., 2007; Diener et al. 2007; Tajčmanová, et al., 2009; Diener & Powell, 2012) there remain considerable challenges and inconsistencies. As discussed in White et al. (2014) the
persistence of margarite-bearing fields in common metapelite compositions represents one such challenge, and it would appear at this stage that an appraisal of how the Na$_2$O–CaO–K$_2$O-bearing phases (mica feldspar and epidote) are interacting in the large systems is required. In addition to uncertainties associated with the $a$–$x$ relations themselves, there are considerable uncertainties regarding oxidation state of several key elements including Mn. To these may be added many geologically-based sources of uncertainty such as in relation to equilibrium volume, including fractionation of components into porphyroblast cores, composition of co-existing fluids and open system behaviour. Given these caveats interpretations that use the results literally, such as to interpret the degree of overstep of reactions based on intersecting compositional isopleths are unlikely to be defensible. Thus, the $a$–$x$ relations presented here and phase diagrams produced from their use are better thought of as a thermodynamic framework in which to interpret metamorphic features rather than a literal and absolute solution. However, these limitations should not prevent the quantitative assessment of metamorphic conditions, but rather should prompt an appropriate degree of uncertainty to be attached to such results.

**Acknowledgements**

### and ### are thanked for their ... reviews. RP acknowledges support from Australian Research Council grants DP0451770 and DP0987731.
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APPENDIX 1: Interaction energies

Following the adoption of micro-$\phi$ for the incorporation of the Mn-end-members in the minerals, as discussed in the text, along with the interaction energies from White et al. (2014), the interaction energies $W_{ij}$ in matrix form are:

\[
\begin{array}{c|cccc}
  \text{g} & \text{py} & \text{gr} & \text{kho} & \text{spss} \\
\hline
  \text{alm} & 2.5 & 5 & 22.6 & 2 \\
  \text{py} & 31 & 5.4 & 2 & \\
  \text{gr} & -15.3 & 0 & \\
  \text{kho} & 29.4 & \\
\end{array}
\]

\[
\begin{array}{c|cccccccc}
  \text{chl} & \text{afchl} & \text{ames} & \text{daph} & \text{ochl1} & \text{ochl4} & \text{f3clin} & \text{mnchl} \\
\hline
  \text{clin} & 17 & 17 & 20 & 30 & 21 & 2 & 15 \\
  \text{afchl} & 16 & 37 & 20 & 4 & 15 & 32 & \\
  \text{ames} & 30 & 29 & 13 & 19 & 26 & \\
  \text{daph} & 18 & 33 & 22 & 10 & \\
  \text{ochl1} & 24 & 28.6 & 25 & \\
  \text{ochl4} & 19 & 31 & \\
  \text{f3clin} & & & & 17 \\
\end{array}
\]

\[
\begin{array}{c|cccccccc}
  \text{bi} & \text{ann} & \text{obi} & \text{east} & \text{tbi} & \text{fbi} & \text{mn} \\
\hline
  \text{phl} & 12 & 4 & 10 & 30 & 8 & 9 \\
  \text{ann} & 8 & 15 & 32 & 13.6 & 6.3 & \\
  \text{obi} & 7 & 24 & 5.6 & 8.1 & \\
  \text{east} & 40 & 1 & 13 & \\
  \text{tbi} & 40 & 30 & \\
  \text{fbi} & & & & 11.6 \\
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APPENDIX 2: Natural assemblage data

A natural assemblage database has been assembled with the focus on mineral assemblages from rocks with a range of MnO content (see Supplementary material). The analyses given in the supplementary material were taken from the following: 1, Grew (1981) — rutile present; 2, Hauzenberger et al. (2001); 3, Sevigny & Ghent (1989); 4, Redler (2011); 5, Kunz (2011); 6, Fraser et al. (2000); 7, Bickle & Archibald (1984); 8, Bosse et al. (2002); 9, Blümel & Schreyer (1977); 10, Droop & Moazzen (2007); 11, Álvarez-Valero et al. (2007) — El Hoyazo locality; 12, Fletcher & Greenwood (1979); 13, Ghent & De Vries (1972); 14, Greenfield (1997); 15, Grew (1981) — rutile absent; 16, Guidotti (1974) — supp. data; 17, Guidotti (1974); 18, Harris (1976); 19, Hodges & Spear (1982); 20, Kamineni (1975); 21, Kawakami et al. (2007); 22, Likhanov et al. (2001); 23, Lui (2004); 24, Mather (1970); 25, Álvarez-Valero et al. (2007) — Mazarrón locality; 26, Mposkos (1989); 27, Novak & Holdaway (1981); 28, Otamendi et al. (1999); 29, Pomroy (2004); 30, Pattison & Vogl (2005); 31, Tiddy (2002); 32, Johnson et al. (2004); 33, Vilà et al. (2007); 34, West et al. (2008); 35, Williams & Grambling, (1990); 36, White (1997); 37, Waters & Whales (1984); 38, Stewart (1942); 39, Thompson et al. (1977); 40, Atherton (1968); 41, Heald (1940); 42, Green (1963); 43, Leake (1958); 44, Chinner (1960); 45, Miyashiroy (1953); 46, Engel & Engel (1960); 47, Phinney (1963); 48, Albee (1965); 49, Hietanen (1956); 50, Chinner (1967); 51, Reinhardt (1968); 52, Davidson & Matheson (1974); 53, Kays & Medaris (1976); 54, Chinner (1962); 55, Herman et al. (1978); 56, Patiño Douce & Beard (1995); 57, Patiño Douce & Beard (1996); 58, Patiño Douce et al. (1993); 59, Vielzeuf & Montel (1997).

A proportion of the data involved wet chemical analyses, with analysed FeO and Fe₂O₃. The data for the remainder of the analyses involve all-Fe-as-FeO. Given the large uncertainties in calculated ferrous-ferric that arise from charge balance calculations (e.g. Powell & Holland, 2008), with the added difficulty of the reliability of stoichiometric constraints for hydrous minerals, the approach followed for analyses involving all-Fe-as-FeO is to adopt heuristics for conversion of FeO to Fe₂O₃: as proportions, cd = 0; ctd = 0.02; g = 0.03; opx = 0.05; chl = 0.1; st = 0.1; and bi = 0.15 (see also the Appendix in White et al., 2014). Given that Mg–Mn exchange reactions are used in the data analysis, the precise values used for the conversion are not important, but it is appropriate to have them in what is considered to be a petrologically consistent order. Charge balance is used to
recalculate the ilmenite analyses.

Processing of the analyses involves mineral calculation in the commonly-used way, on a specified number of oxygens, and assuming the full complement of hydroxyls for the hydrous minerals (except for biotite as a consequence of the oxy-substitution used for Ti). The cations are then used to calculate the composition parameters as defined in the Appendix of White et al. (2014). The order parameters, $Q$, for the Fe–Mg order-disorder in the minerals (and the Mg–Al order-disorder in chlorite) are calculated by solving the appropriate internal equilibria in each mineral using the $a$-$x$ relations given in Appendix 1, and the $\Delta H$ of the internal equilibria given in the Appendix of White et al. (2014).
The Mg-Mn exchange reactions can be written in a standard form, involving one cation exchange between minerals A and B, with A–B for Mg–Mn meaning

\[ \text{Mn, } \text{A} + \text{Mg, } \text{B} = \text{Mg, } \text{A} + \text{Mn, } \text{B} \]

or as an example, g–bi for Mg–Mn meaning that we are considering the equilibrium involving

\[ \frac{1}{3}\text{spss, g} + \frac{1}{3}\text{phl, bi} = \frac{1}{3}\text{py, g} + \frac{1}{3}\text{mn, bi} \]

The thermodynamics, in the form \( \Delta h = \Delta G^\circ + RT \ln K \), are evaluated for each such reaction for each appropriate natural assemblage mineral pair. In this, \( \Delta G^\circ \) is calculated using Holland & Powell (2011), and \( K \) is calculated from the compositional and order parameters calculated from the mineral compositions as outlined in Appendix 2. \( \Delta h \) can be thought of as a \( \Delta \text{DQF} \) on a one cation exchange basis. Given that the enthalpies of the magnesian end-members are taken to be well-known, \( \Delta h \) relates just to the manganese end-members, e.g. for the above example.

\[ \Delta h_{\text{g,bi}} = -\frac{1}{3}\text{DQF}_{\text{spss, g}} + \frac{1}{3}\text{DQF}_{\text{mn,bi}} \]

As noted above, the properties of the pyrophanite and spessartine end-members in the Holland & Powell (2011) dataset are not based on Mahar et al. (1997). Therefore, in principle, this means that

\[ \Delta h_{\text{g,ilm}} = -\frac{1}{3}\text{DQF}_{\text{spss, g}} + \text{DQF}_{\text{ptt,ilm}} = 0 \]

This can be assessed with the database here, and Fig. A3-1a,b shows that this is the case within error (49 data points). In Fig. A3-2, a selection of \( \Delta h \) plots show the nature of the data, the dotted line being the median of the data, the band representing its uncertainty (see below), and the solid line the result of the least squares analysis of all of the data below.

In the following table, \( n \) is the number of mineral pairs involved. \( \Delta h \) is the median of the natural assemblage values; \( \sigma_{\Delta h}^{\text{distr}} \) is an estimate of the standard deviation on this value using the normalised median absolute deviation, nMAD (e.g. Powell et al., 2002). Medians are used to try and downplay the effect of the scatter and outliers in the data. Treating
\( \sigma_{\Delta h}^{\text{distr}} \) now as a Gaussian estimate of standard deviation on the distribution of the data, this divided by the square root of the number of data gives an estimate of the standard deviation on the mean of \( \Delta h \). This is what would then used to represent the data in the next stage of the data, if the data are homoscedastic, in other words if they are all uncorrelated. But the data are most likely correlated given that they are not individually from separate studies. As a gross approximation to account for this, we use \( \sigma_{\Delta h} = 2\sigma_{\Delta h}^{\text{mean}} \), the last column of the following table. That this is appropriate is established \emph{a posteriori} below. the uncertainty bands in Figs A3-1, A3-2 are for \( 2\sigma_{\Delta h} \)

<table>
<thead>
<tr>
<th>A</th>
<th>B</th>
<th>n</th>
<th>( \Delta h )</th>
<th>( \sigma_{\Delta h}^{\text{distr}} )</th>
<th>( 2\sigma_{\Delta h}^{\text{mean}} )</th>
</tr>
</thead>
<tbody>
<tr>
<td>g</td>
<td>bi</td>
<td>149</td>
<td>2.623</td>
<td>3.821</td>
<td>0.626</td>
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<td>6.595</td>
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<td>cd</td>
<td>27</td>
<td>4.815</td>
<td>9.239</td>
<td>3.556</td>
</tr>
<tr>
<td>g</td>
<td>ctd</td>
<td>25</td>
<td>1.298</td>
<td>4.340</td>
<td>1.736</td>
</tr>
<tr>
<td>g</td>
<td>st</td>
<td>47</td>
<td>-0.325</td>
<td>1.661</td>
<td>0.484</td>
</tr>
<tr>
<td>g</td>
<td>opx</td>
<td>23</td>
<td>-0.964</td>
<td>4.814</td>
<td>2.008</td>
</tr>
<tr>
<td>bi</td>
<td>chl</td>
<td>29</td>
<td>0.490</td>
<td>2.803</td>
<td>1.041</td>
</tr>
<tr>
<td>bi</td>
<td>cd</td>
<td>64</td>
<td>-1.639</td>
<td>3.132</td>
<td>0.783</td>
</tr>
<tr>
<td>bi</td>
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<td>3.308</td>
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<tr>
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<td>5.661</td>
<td>1.570</td>
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<tr>
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<td>0.765</td>
<td>1.476</td>
<td>1.044</td>
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<td>chl</td>
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<td>1.685</td>
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<tr>
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<td>st</td>
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<td>1.084</td>
<td>0.654</td>
</tr>
<tr>
<td>cd</td>
<td>opx</td>
<td>13</td>
<td>-8.310</td>
<td>4.593</td>
<td>2.548</td>
</tr>
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<td>ctd</td>
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<td>14</td>
<td>0.180</td>
<td>3.181</td>
<td>1.701</td>
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</table>

Determining the “best” \( \Delta h \) values from this table is a weighted least squares problem. Assuming that there is no enthalpy modification needed for spessartine, the analysis gives the enthalpy modifications for the individual manganese end-members, on a one cation
Multiplied by the number of Mn in the end-member formulae, this gives the \( dqf \) of the manganese end-members.

\[
\begin{array}{cccccc}
\text{bi} & \text{chl} & \text{cd} & \text{ctd} & \text{st} & \text{opx} \\
\hline
h & -2.63 & -2.61 & -2.10 & 0.66 & -0.04 & 3.34 \\
\end{array}
\]

The \( \sigma_{fit} \) of the least squares is 1.24, reflecting that the \( \sigma_h \) used were not inappropriate. The correlation coefficient matrix of the \( dqf \) is:

\[
\begin{array}{cccccc}
\text{bi} & \text{chl} & \text{cd} & \text{ctd} & \text{st} & \text{opx} \\
\hline
\text{dqf} & -7.89 & -13.03 & -4.21 & 0.66 & -0.17 & 6.68 \\
\end{array}
\]

\[
\begin{array}{cccccc}
\rho & \text{bi} & \text{chl} & \text{cd} & \text{ctd} & \text{st} & \text{opx} \\
\hline
\text{bi} & 1 & 0.429 & 0.526 & 0.131 & 0.337 & 0.275 \\
\text{chl} & 0.429 & 1 & 0.445 & 0.267 & 0.508 & 0.144 \\
\text{cd} & 0.526 & 0.445 & 1 & 0.126 & 0.275 & 0.232 \\
\text{ctd} & 0.131 & 0.267 & 0.126 & 1 & 0.237 & 0.043 \\
\text{st} & 0.337 & 0.508 & 0.275 & 0.237 & 1 & 0.105 \\
\text{opx} & 0.275 & 0.144 & 0.232 & 0.043 & 0.105 & 1 \\
\end{array}
\]
Casting the least squares result in the original form, we get:

\[
\begin{array}{cccccccc}
A & B & \Delta h^\text{obs} & 2\sigma_{h}\text{obs} & \Delta h^\text{calc} & 2\sigma_{h}\text{calc} & e & e^* \\
\hline
g & bi & 2.62 & 1.25 & 2.63 & 1.45 & 0.01 & 0.01 \\
g & chl & 3.58 & 5.62 & 2.61 & 1.76 & -0.98 & -0.35 \\
g & cd & 4.82 & 7.11 & 2.10 & 2.19 & -2.71 & -0.76 \\
g & ctd & 1.30 & 3.47 & -0.66 & 3.05 & -1.96 & -1.13 \\
g & st & -0.32 & 0.97 & 0.04 & 1.24 & 0.37 & 0.76 \\
g & opx & -0.96 & 4.02 & -3.34 & 3.44 & -2.38 & -1.18 \\
b & chl & 0.49 & 2.08 & -0.02 & 1.74 & -0.51 & -0.49 \\
b & cd & -1.64 & 1.57 & -0.53 & 1.88 & 1.11 & 1.42 \\
b & st & -1.09 & 2.02 & -2.59 & 1.56 & -1.50 & -1.49 \\
b & opx & -6.34 & 3.14 & -5.97 & 3.35 & 0.37 & 0.23 \\
chl & cd & 0.76 & 2.09 & -0.50 & 2.11 & -1.27 & -1.21 \\
chl & ctd & -5.63 & 3.37 & -3.27 & 3.09 & 2.36 & 1.40 \\
chl & st & -2.45 & 1.31 & -2.56 & 1.56 & -0.12 & -0.18 \\
chl & opx & -8.31 & 5.10 & -5.45 & 3.63 & 2.86 & 1.12 \\
ctd & st & 0.18 & 3.40 & 0.71 & 3.01 & 0.53 & 0.31 \\
\end{array}
\]

This shows that the fit of the data is good, with \( e = \Delta h^\text{calc} - \Delta h^\text{obs} \) the residuals, and \( e^* \) the residuals normalised to the original specified uncertainties on the data, \( \sigma_{h}\text{obs} \).
**Figure captions**

**Fig. 1:** Petrogenetic grids in the MnKFMASH system for subsolidus and suprasolidus conditions. The square inset shows the in-excess phases used in the different parts of the diagram. The horizontal ticks on select univariants give the value of m(g) of garnet along the reaction. The set of open triangles at 4.5 kbar indicate the \( P-T \) conditions for the AFM compatibility triangles in Fig. 2. The filled triangles at 12 and 6 kbar show the conditions for the MnO–FeO–MgO compatibility triangles in Fig. 3.

**Fig. 2:** Al\(_2\)O\(_3\)–FeO–MgO (AFM) compatibility triangles for a sequence of temperatures at 4.5 kbar. The MnKFMASH system is reduced to AFM by taking garnet, muscovite, quartz and H\(_2\)O to be in excess. The diagrams show the changing divariant to trivariant phase relationships on crossing reactions 3 to 6 (see text for details). The \( P-T \) conditions for each compatibility triangle are given on the figure and additionally shown as a series of open triangles in Fig. 1.

**Fig. 3:** MnO–FeO–MgO compatibility triangles for conditions either side of reaction 3 at 12 kbar (Fig. 3a, b) and 6 kbar (Fig. 3c, d). The MnKFMASH system is reduced to MnO–FeO–MgO by taking kyanite, muscovite, quartz and H\(_2\)O to be in excess and is thus only applicable to aluminous metapelites. The \( P-T \) conditions for each compatibility triangle are given on the figure and additionally shown as a series of filled triangles in Fig. 1.

**Fig. 4:** MnKFMASH pseudosections constructed for three different bulk compositions in terms of the A/AFM and \( x_{Fe} \) proportions, which are given on each pseudosection. Each of the pseudosections is calculated for a MnO content of 0.1 mol % and additional garnet-in lines are shown for an MnO content of 0.05 mol % (red dashed line) and for the MnO-free system (solid red line labelled MnO = 0). (a) \( P-T \) pseudosection from 0.4 to 12 kbar and 450–950 °C. For the subsolidus calculations, H\(_2\)O was taken to be in excess, the H\(_2\)O content for the suprasolidus calculations was set such that the assemblage at the solidus was just saturated in H\(_2\)O (see Table 1). The bulk rock composition is that from fig. 4 in White *et al.* (2014). (b) \( P-T \) pseudosection from 0.4 to 12 kbar and 450–700 °C for a more magnesian composition.
than (Fig. 4a) but with the same Al₂O₃ content. (c). P–T pseudosection for a more aluminous composition but with the same $x_{Fe}$ as Fig. 4b.

**Fig. 5:** MnNCKFMASHTO P–T pseudosection for subsolidus conditions based on the synthetic metapelite composition from Fig. 5b in White et al. (2014) but with 0.1 mol% MnO. The garnet-in boundary is shown as a thick red line. Garnet-in boundaries for bulk MnO contents of 0 mol% (thin red line), 0.05 mol% (thin dashed red line) and 0.15 mol% (thin dotted red line) are additionally shown. Several zero mode boundaries are highlighted in colour (see legend for explanation).

**Fig. 6:** MnNCKFMASHTO P–T pseudosection for subsolidus conditions for a synthetic aluminous metapelite composition. The bulk composition used is that from Fig. 7 in White et al. (2014) with 0.1 mol% added. The garnet-in boundary is shown as a thick red line. Garnet-in boundaries for bulk MnO contents of 0 mol% (thin red line), 0.05 mol% (thin dashed red line) and 0.15 mol% (thin dotted red line) are additionally shown. Several zero mode boundaries are highlighted in colour (see legend for explanation).

**Fig. 7:** MnNCKFMASHTO P–T pseudosection for subsolidus conditions calculated for the average greenschist facies metapelite composition from Ague (1991). The CaO content was reduced slightly to account for the likely presence of apatite in most metapelites. The value for O, representing the Fe₂O₃ component was set at 0.498 mol% to give a Fe₂O₃/(FeO + Fe₂O₃) value of 0.077. The bulk composition used has 0.284 mol% MnO and garnet is stable throughout much of the diagram. Garnet-in lines are also shown for bulk MnO contents of 0.05 and 0.15 mol%. Several zero mode boundaries are highlighted in colour (see legend for explanation).

**Fig. 8:** MnNCKFMASHTO P–T pseudosection for subsolidus conditions calculated for the average amphibolite facies metapelite composition from Ague (1991). The value for O, representing the Fe₂O₃ component was set at 0.602 mol% to give a Fe₂O₃/(FeO + Fe₂O₃) value of 0.081. The bulk composition used has 0.175 mol% MnO and garnet has a smaller stability range than for Fig. 7. Garnet-in lines are also shown for bulk MnO contents of 0.05 and 0.15 mol%. Several zero mode boundaries are highlighted in colour (see legend for explanation).
**Fig. 9:** MnNCKFAMSHTO $P$–$T$ pseudosection for suprasolidus conditions calculated for the same composition as Fig. 8. The $H_2O$ content was set such that the solidus was just fluid saturated at close to 9 kbar. The garnet-in line for the MnO-free system is also shown as a thin red line labelled MnO = 0. Several zero mode boundaries are highlighted in colour (see legend for explanation).

**Fig. 10:** A $T$–$x_{\text{MnO}}$ (Fig. 10a) and $P$–$x_{\text{MnO}}$ (Fig. 10b) pseudosection based on the greenschist facies metapelite composition presented in Ague (1991). For both diagrams the $x$ axis ranges from MnO = 0 mol % at $x = 0$ to MnO = 0.3 mol % at $x = 1$. In both diagrams the garnet in line is shown as a thick red line and contours of garnet mode are shown as thin red lines. (a). A $T$–$x_{\text{MnO}}$ pseudosection from 400°C to the wet solidus calculated for a pressure of 6 kbar. (b). A $P$–$x_{\text{MnO}}$ pseudosection for 0.4–10 kbar, calculated for a temperature of 580°C.

**Fig. A3-1:** Plots of $\Delta h$ versus manganese compositional parameters ($m$) for garnet–ilmenite pairs. Natural assemblage data, ferric iron not analysed (circles); natural assemblage data, wet chemistry (squares); experimental data, ferric iron not analysed (diamonds). (see text).

**Fig. A3-2:** Plots of $\Delta h$ versus manganese compositional parameters ($m$). (a–f) select mineral pairs involving garnet with biotite, chloritoid and staurolite. (g–l) select mineral pairs involving biotite with chlorite, cordierite and orthopyroxene. Natural assemblage data, ferric not analysed (circles); natural assemblage data, wet chemistry (squares); experimental data, ferric not analysed (diamonds). (see text).
Table 1: Bulk rock compositions used in the construction of pseudosections

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<th>mol. %</th>
<th>H₂O</th>
<th>SiO₂</th>
<th>Al₂O₃</th>
<th>CaO</th>
<th>MgO</th>
<th>FeO</th>
<th>K₂O</th>
<th>Na₂O</th>
<th>TiO₂</th>
<th>MnO</th>
<th>O</th>
</tr>
</thead>
<tbody>
<tr>
<td>Fig. 4a</td>
<td>6.553*</td>
<td>68.691</td>
<td>9.860</td>
<td>—</td>
<td>4.006</td>
<td>7.632</td>
<td>3.157</td>
<td>—</td>
<td>0.100</td>
<td>—</td>
<td>—</td>
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<tr>
<td>Fig. 4b</td>
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<td>—</td>
<td>5.480</td>
<td>6.975</td>
<td>3.378</td>
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<td>0.107</td>
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<td>Fig. 4c</td>
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+, H₂O in excess; *, H₂O taken as in-excess for subsolidus part of diagram
Figure 2:
Figure 3:

(a) MgO

(b) MgO

(c) MgO

(d) MgO

MgO, g, ct, chl, MnO, FeO, +ky, +mu, +q, +H2O

501 °C, 6 kbar

562 °C, 12 kbar

502 °C, 12 kbar

504 °C, 12 kbar
Figure 6:
Figure 9: MnNCKFMASHTO
White et al fig. 3-1
Figure 3-2: Diagrams showing the relationship between ∆h and m(g) for different labels:

(a) g-bi
(b) g-bi
(c) g-ctd
d) g-ctd
(e) g-st
(f) g-st
White et al. fig. 3-2 - cont.